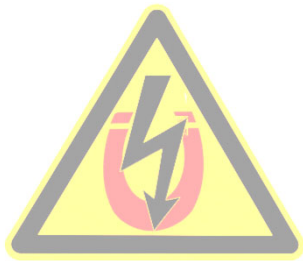


PHYS 301

Electricity and Magnetism



Dr. Gregory W. Clark
Fall 2019

Today!

- Faraday's Law
- Maxwell's Law
- Maxwell's Equations

Faraday's Law

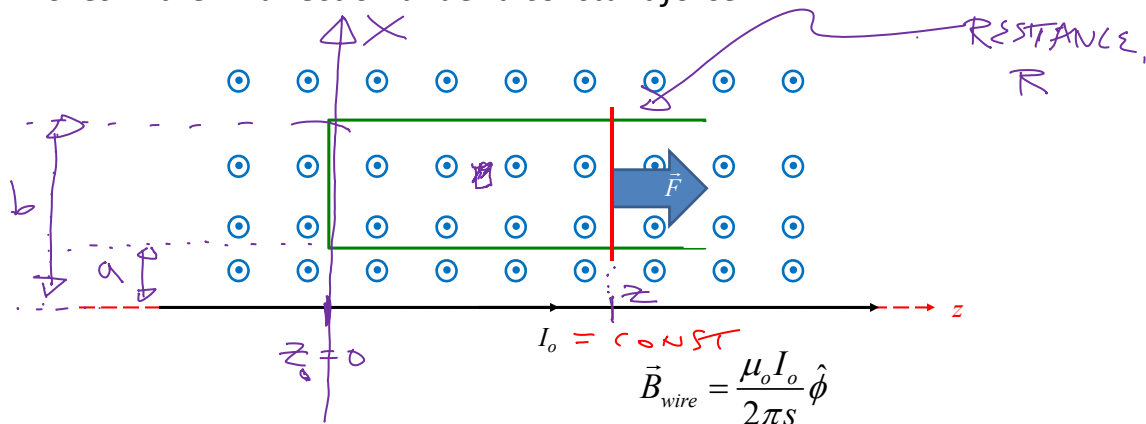
$$\oint \vec{E} \cdot d\vec{\ell} = -\frac{d\Phi_B}{dt}$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Faraday's Law

Current induced in loop near current carrying wire

- Find induced current in green wire when red conducting "slider" moves in the $+z$ direction under a constant force.




Faraday's Law

QUASISTATIC APPROXIMATION:

Changing \vec{B} 's produce \vec{E} 's, but often
the formalism of magnetostatics
is used to (approximately) describe
the changing \vec{B} 's!

[~ okay, if \vec{B} doesn't change "too rapidly!"]


e.g., fast enough to produce EM waves!

Faraday's Law

➤ THINGS TO KEEP IN MIND ...

MOTIONAL EMFs - MAGNETIC FORCE SETS UP THE EMF, BUT DOES NOT TO WORK!

(it can't – why?!)

WHATEVER IS MOVING THE WIRES (LOOP – THIS FORCE IS TRANSMITTED
TO CHARGES VIA STRUCTURE OF THE WIRE.

MAGNETIC INDUCTION – MAGNET MOVES (OR FIELD CHANGES) BUT LOOP IS STATIONARY;
∴ CHARGES ARE NOT MOVING SO NO MAGNETIC FORCE IS INVOLVED!

FORCE IS DUE TO AN ELECTRIC FIELD PRODUCED BY THE CHANGING MAGNETIC FIELD:

$$\vec{\nabla} \times \vec{E} = -\partial \vec{B} / \partial t$$

❖ Fascinating that the two phenomena can be described in the same way.
As Einstein showed, it's the relative motion that is important.

MAXWELL'S EQUATIONS

So far, we have . . .

$$\begin{aligned}\vec{\nabla} \cdot \vec{E} &= \rho / \epsilon_0 & \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} & \vec{\nabla} \times \vec{B} &= \mu_0 \vec{J}\end{aligned}$$

★ Problem: $\vec{\nabla} \cdot (\vec{\nabla} \times \vec{V}) = 0 \quad \forall \text{ vectors } \vec{V}$

Consider Faraday's law . . .

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{E}) = -\frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{B}) = 0 \quad \text{Okay for both terms!}$$

MAXWELL'S EQUATIONS

★ Problem: $\vec{\nabla} \cdot (\vec{\nabla} \times \vec{V}) = 0 \quad \forall \text{ vectors } \vec{V}$

Consider Ampere's law . . .

$$\underbrace{\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B})}_{=0} = \underbrace{\mu_0 (\vec{\nabla} \cdot \vec{J})}_{\neq 0, \text{ in general!}}$$

For steady currents,

$$\vec{\nabla} \cdot \vec{J} = 0 = -\frac{\partial \rho}{\partial t}$$

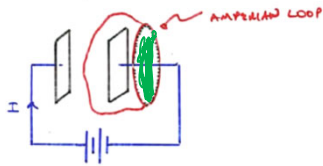
and it works . . .

But beyond this realm of **magnetostatics** it does not hold!

MAXWELL'S EQUATIONS

• The standard example:

Apply Ampere's law to the charging of a capacitor



$$\oint \vec{B} \cdot d\vec{\ell} = \mu_o I_{ENCL} = \mu_o \int_{surf} \vec{J} \cdot d\vec{A}$$

★ Look at I_{ENCL} through . . .

➤ Surface in plane of **LOOP**:

$$I_{ENCL} = I$$

➤ Surface bounded by **LOOP**, but which passes between plates:

$$I_{ENCL} = 0$$



*This was not a problem in magnetostatics
because the currents were steady
Here, charge is piling up somewhere!!
(where?!)*

MAXWELL'S EQUATIONS

★ MAXWELL'S FIX:

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) = 0 = \mu_o (\vec{\nabla} \cdot \vec{J}) + \underline{\text{something!}}$$

$$\text{since } \vec{\nabla} \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$$

$$\text{and } \rho = \epsilon_o \vec{\nabla} \cdot \vec{E}$$

$$\text{then } \vec{\nabla} \cdot \vec{J} = -\epsilon_o \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{E}) = -\vec{\nabla} \cdot \left[\epsilon_o \frac{\partial \vec{E}}{\partial t} \right]$$

$$\text{So, let the } \underline{\text{something}} \text{ be } -\mu_o \vec{\nabla} \cdot \vec{J} \text{ written as } + \vec{\nabla} \cdot \left[\underline{\mu_o \epsilon_o \frac{\partial \vec{E}}{\partial t}} \right]$$

MAXWELL'S EQUATIONS

THEN WE HAVE $\underline{\underline{\vec{\nabla} \times \vec{B} = \mu_o \vec{J} + \mu_o \epsilon_o \frac{\partial \vec{E}}{\partial t}}}$

- ❖ Causes no trouble in magnetostatics,
where $\vec{E} = \text{constant}$
- ❖ Implies that changing $\vec{E} \Rightarrow \vec{B}$
Just like changing $\vec{B} \Rightarrow \vec{E}$ (*Faraday*)
- ❖ Crucial in propagation of waves
- ❖ Maxwell: $\vec{J}_d \equiv \epsilon_o \frac{\partial \vec{E}}{\partial t} = \text{DISPLACEMENT "CURRENT"}$

MAXWELL'S EQUATIONS IN FREE SPACE

$$\begin{aligned} \vec{\nabla} \cdot \vec{E} &= \rho / \epsilon_o & \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} & \vec{\nabla} \times \vec{B} &= \mu_o \vec{J} + \mu_o \epsilon_o \frac{\partial \vec{E}}{\partial t} \\ \vec{F} &= q(\vec{E} + \vec{v} \times \vec{B}) \end{aligned}$$

MAXWELL'S EQUATIONS
IN MATTER

$$\vec{\nabla} \cdot \vec{D} = \rho_f \qquad \vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \qquad \vec{\nabla} \times \vec{H} = \vec{J}_f + \frac{\partial \vec{D}}{\partial t}$$

where $\vec{D} = \epsilon_o \vec{E} + \vec{P}$

$$\vec{H} = \frac{1}{\mu_o} \vec{B} - \vec{M}$$

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$